

CTEQ GLOBAL QCD ANALYSIS

- * PDFs for showering programs
- * PDFs with a variable QCD coupling

Pavel Nadolsky

Southern Methodist University
Dallas, TX, U.S.A.

in collaboration with

J. Huston, M. Guzzi, H.-L. Lai, Z. Li, S. Mrenna, J. Pumplin,
D. Stump, W.-K. Tung, and C.-P. Yuan

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CTEQ-Tung Et Al.: ongoing activities

- PDFs for leading-order Monte-Carlo programs (*arXiv:0910.4183*)
- Uncertainty in α_s in the CTEQ PDF analysis (*arXiv:1004.4624*)
- General-purpose NLO PDF fits
 - ▶ CTEQ6.6 set (published in 2008) → CT09 (not released)
→ CT10 (pre-released within CTEQ)
 - ▶ combined HERA and Tevatron Run-2 lepton asymmetry data are included
 - ▶ new statistical methods and parametrization forms
- Benchmarking of heavy-quark contributions
2009 Les Houches Proceedings
- Constraints on color-octet fermions (*with Berger, Guzzi, Olness*)
- Exploration of statistical methods and PDF parametrizations
(*Pumplin, arXiv:0909.0268 and 0909.5176*)

CT09MC PDF's for leading-order

Monte-Carlo generators at the LHC

Lai, Huston, Mrenna, P. N., Stump, Tung, Yuan, JHEP 1004, 035 (2010)

Related studies: Sherstnev, Thorne, arXiv:0711.2473; Jung, Sjostrand @ the HERA-LHC workshop

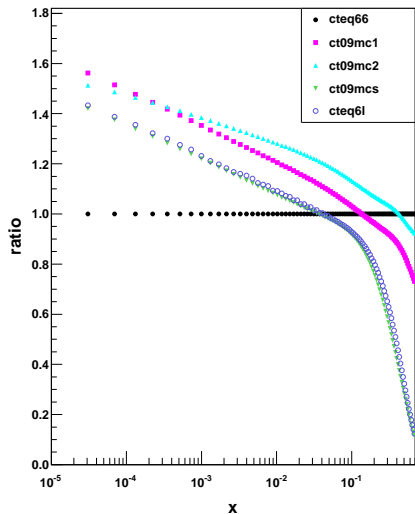
- **Leading-log** showering programs (PYTHIA, HERWIG,...) will remain ubiquitous tools in the observable future
- We wish to construct improved PDFs for showering programs **at the LHC**, to be used when NLO calculations (with or without showering) are not applicable for some reason
- These PDFs do not have to be “strictly at leading order”, because showering introduces contributions beyond LO

Requirements for LO-MC PDFs

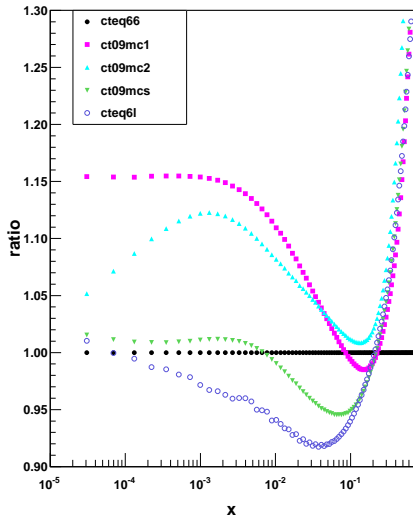
- **LO** DGLAP evolution
- Steep **LO** gluon PDF $g(x, Q)$ at small x , required by models of underlying event
- when combined with LO matrix elements, **reproduce NLO rapidity distributions for key LHC processes** (production of $W^\pm, Z^0, H^0 \dots$)
- \Rightarrow are about the same as NLO PDFs at large x , or exceed them

Resulting PDFs at $Q = 85 \text{ GeV}$

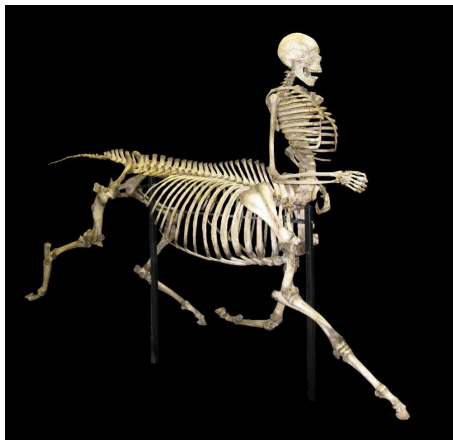
Q value: 85 GeV, for parton: g



Q value: 85 GeV, for parton: u



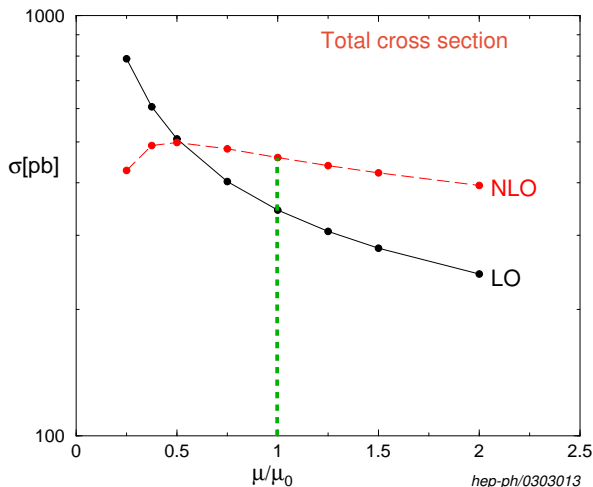
Is this liminal creature viable?



Let's examine the bones, focusing on fixed-order cross sections.

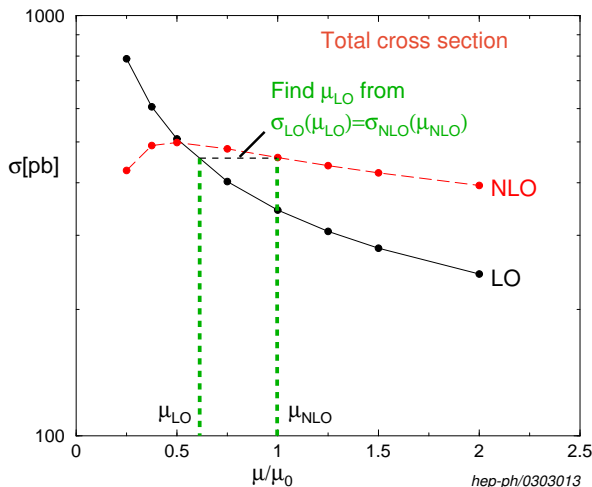
Modifications by parton showering (in PYTHIA) are smaller than the dependence on PDFs (\Rightarrow backup slides)

Approximating NLO cross sections by LO cross sections



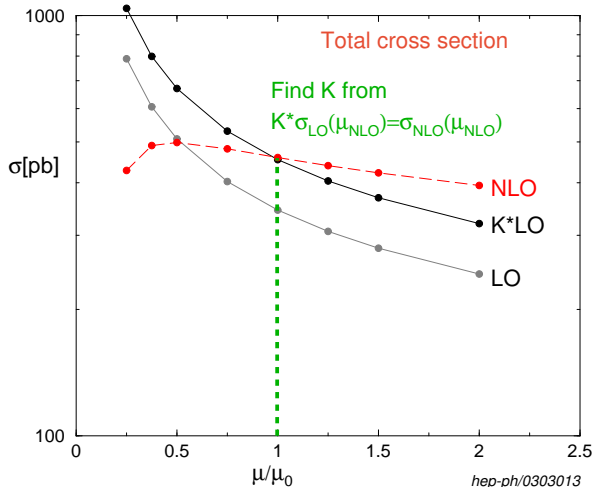
- NLO cross sections: more predictive
- LO cross sections: more uncertain (=flexible)

Approximating NLO cross sections by LO cross sections



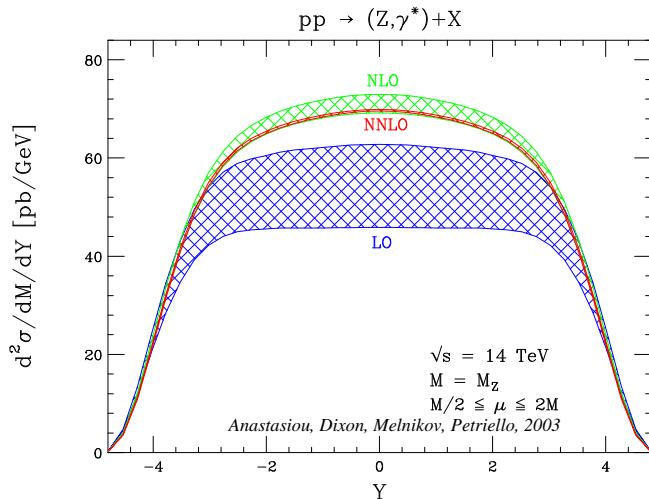
Option 1: Optimize the factorization scale μ_{LO}

Approximating NLO cross sections by LO cross sections



Option 2: Optimize the factor $K = \sigma_{NLO}(\mu_{NLO})/\sigma_{LO}(\mu_{NLO})$ (i.e., the floating normalization of σ_{LO})

Approximating NLO cross sections by LO cross sections



For $d\sigma_{NLO}/dy$, etc., optimize **both** μ_{LO} and K to reproduce normalization and shape

CT09MC PDFs: selection of data

■ Perform a **LO** fit to

- ▶ CTEQ6.6 set of experimental data from DIS, production of vector bosons and jets (2700 points)
- ▶ **pseudodata** at $\sqrt{s} = 14$ TeV, containing **NLO predictions** probing typical combinations of PDFs at various (x, Q) :

$$\diamond d\sigma/dy \text{ for } pp \rightarrow W^\pm X, pp \rightarrow Z^0 X, pp \rightarrow H^0 X; d\sigma/dM_{t\bar{t}} \text{ for } pp \rightarrow t\bar{t} X; d\sigma/dM_{b\bar{b}} \text{ for } gg \rightarrow b\bar{b} \text{ at } M_{b\bar{b}} \text{ of } 10\text{-}50 \text{ GeV}$$

All cross sections are at fixed order; modifications by showering are examined outside of the fit

Vary K_i and/or $\mu_{LO,i}$ in the fit

Three CT09MC PDF sets

PDF set	CT09MCS	CT09MC1	CT09MC2
$\mu_{LO,i}$	Varied	Fixed	Fixed
K_i	Varied	Varied	Varied
Loops in α_s	2	1	2
Momentum sum rule	enforced	relaxed	relaxed

- relaxation of the momentum sum rule (*Sherstnev, Thorne*) by 10-15% produces a more flexible $g(x, \mu)$, smaller $K_i - 1$

Reconciling LO with the global data

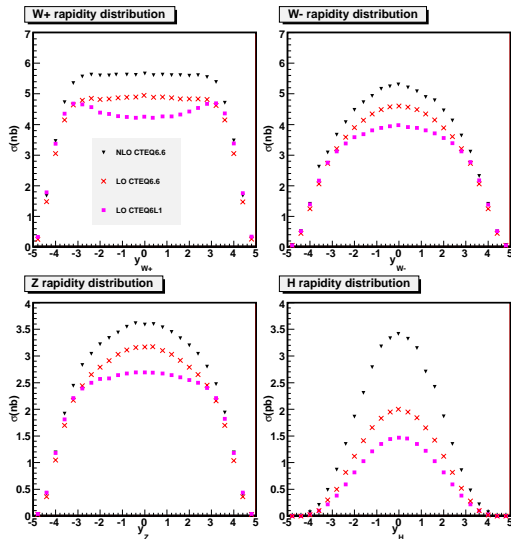
- $\chi^2/d.o.f.$ is worse at LO than at NLO by 20-30%
- When the (mostly high- Q) LHC pseudodata are included, the agreement with the (mostly low- Q) real data deteriorates

Therefore:

- pseudodata errors (χ^2 weights) are chosen to balance between the agreement with the real data and pseudodata
- two-loop α_s improves χ^2
- Variable $\mu_{LO,i}$ and K_i improve χ^2
- relaxed momentum sum **slightly** improves χ^2

$d\sigma/dy$ for conventional LO PDFs

LO ME-LO PDF (CTEQ6L1) vs. LO-NLO (CTEQ6.6M) vs. NLO-NLO (CTEQ6.6M)

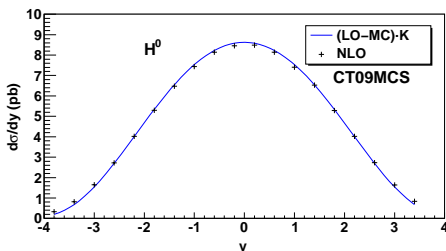
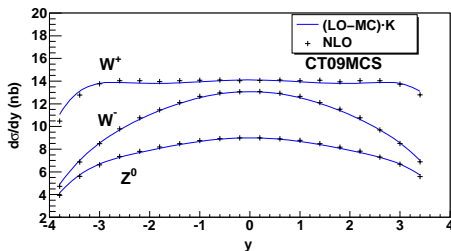


- Significant differences in normalizations and shapes between LO-CTEQ6L1 and LO-CTEQ6.6M predictions
- Unphysical forward-backward peaking in the LO-CTEQ6L1 $d\sigma/dy$ for W^+ production

CT09MCS (scale) set

Two-loop $\alpha_s(\mu)$; vary μ and K ; enforce momentum sum

Agreement with normalization and shape of NLO cross sections

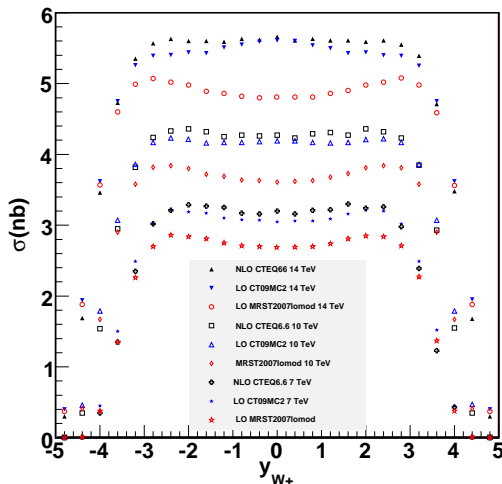


	W^+	W^-	Z	H	$t\bar{t}$	$b'\bar{b}'$
μ_i	$1.96 M_W$	$1.96 M_W$	$1.96 M_Z$	$1.06 M_H$	$1.41 M_t$	$0.40 M_{b'\bar{b}'}$
K_i	1.11	1.09	1.09	1.87	2.09	4.09

CT09MC1 and 2 sets

One- and two-loop $\alpha_s(\mu)$; vary μ and K ; relax momentum sum

W+ rapidity distribution



Best-fit K factors

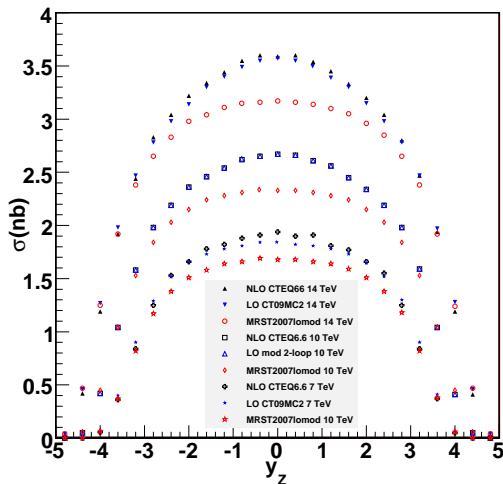
	MC1	MC2
W^+	1.00	1.02
W^-	0.99	1.00
Z^0	0.98	1.00
H^0	1.22	1.32
$t\bar{t}$	1.09	1.09
$b'\bar{b}'$	2.70	3.13
mom. sum	1.10	1.14

LO cross sections are plotted without K factors

CT09MC1 and 2 sets

One- and two-loop $\alpha_s(\mu)$; vary μ and K ; relax momentum sum

Z rapidity distribution



Best-fit K factors

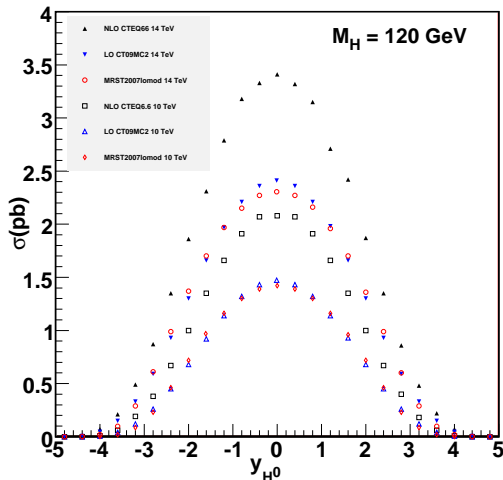
	MC1	MC2
W^+	1.00	1.02
W^-	0.99	1.00
Z^0	0.98	1.00
H^0	1.22	1.32
$t\bar{t}$	1.09	1.09
$b'\bar{b}'$	2.70	3.13
mom. sum	1.10	1.14

LO cross sections are plotted without K factors

CT09MC1 and 2 sets

One- and two-loop $\alpha_s(\mu)$; vary μ and K ; relax momentum sum

SM Higgs boson rapidity distribution



Best-fit K factors

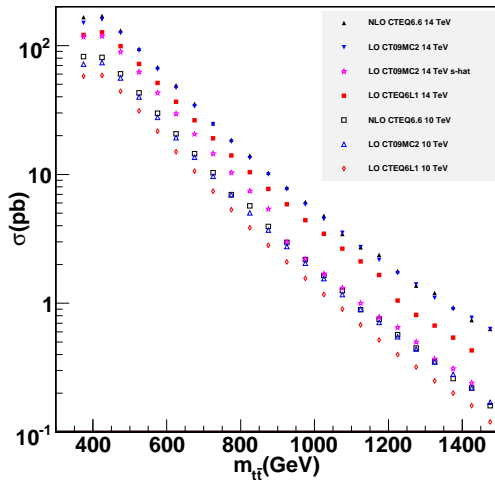
	MC1	MC2
W^+	1.00	1.02
W^-	0.99	1.00
Z^0	0.98	1.00
H^0	1.22	1.32
$t\bar{t}$	1.09	1.09
$b'\bar{b}'$	2.70	3.13
mom. sum	1.10	1.14

LO cross sections are plotted without K factors

CT09MC1 and 2 sets

One- and two-loop $\alpha_s(\mu)$; vary μ and K ; relax momentum sum

$t\bar{t}$ mass distribution



Best-fit K factors

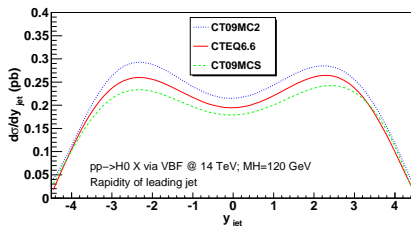
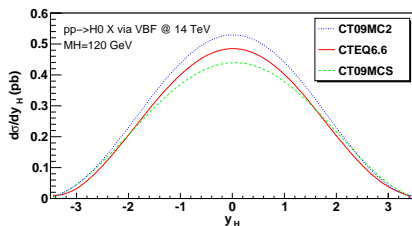
	MC1	MC2
W^+	1.00	1.02
W^-	0.99	1.00
Z^0	0.98	1.00
H^0	1.22	1.32
$t\bar{t}$	1.09	1.09
$b'\bar{b}'$	2.70	3.13
mom. sum	1.10	1.14

LO cross sections are plotted without K factors

CT09MC approximation for other LHC cross sections...

... may or may be not adequate in general, depending on the similarity to the fitted cross sections

...works reasonably well for $pp \rightarrow H^0 X$ via vector boson fusion, for a few other examined processes



Uncertainty in α_s in the CTEQ6.6 and CT10 PDF analysis

arXiv:1004.4624

- Two leading theoretical uncertainties in LHC processes are due to α_s and the PDFs
- These are not independent uncertainties; how can one quantify their correlation?
- Which central $\alpha_s(M_Z)$ and which error on $\alpha_s(M_Z)$ are to be used with the existing PDFs?
- What are the consequences for key LHC processes ($gg \rightarrow H^0$, etc.)?

Uncertainty in α_s in the CTEQ-TEA PDF analysis

arXiv:1004.4624

Recent activity to examine these questions, e.g.:

■ MSTW (arXiv:0905.3531)

- ▶ $\alpha_s(M_Z)$ is an **output** of the global fit (constrained by the hadronic scattering only)
- ▶ several sets of error PDFs, each with its own $\alpha_s(M_Z)$ value \Rightarrow lengthier calculations
- ▶ The α_s uncertainty and PDF uncertainty are inseparable

■ NNPDF (in 2009 Les Houches Proceedings, arXiv:1004.0962):

- ▶ $\alpha_s(M_Z) = 0.119 \pm 0.002$ is taken as an **input**
- ▶ α_s -PDF correlation is examined with ~ 1000 PDF replicas and found to be small

■ H1+ZEUS (arXiv:0911.0884): sensitivity of the HERAPDF set to $\delta\alpha_s(M_Z) = \pm 0.002$ is explored

CTEQ6.6FAS analysis

- Take the “world-average” $\alpha_s(M_Z) = 0.118 \pm 0.002$ as an **input**:

$$\alpha_s(M_Z)|_{\text{in}} = 0.118 \pm 0.002 \text{ at } 90\% \text{ C.L.}$$

- Find the theory parameter $\alpha_s(M_Z)$ as an **output** of a global fit (CTEQ6.6FAS):

$$\alpha_s(M_Z)|_{\text{out}} = 0.118 \pm 0.0019 \text{ at } 90\% \text{ C.L.}$$

- The combined PDF+ α_s uncertainty is estimated as

$$\Delta X = \frac{1}{2} \sqrt{\sum_{i=1}^{22+1} \left(X_i^{(+)} - X_i^{(-)} \right)^2}$$

- Problem:** each PDF set comes with its own $\alpha_s \Rightarrow$ cumbersome
- A simple workaround exists!**

A quadrature sum reproduces the α_s -PDF correlation

H.-L. Lai, J. Pumplin

Theorem

In the quadratic approximation, the total α_s +PDF uncertainty $\Delta\sigma$ of the CTEQ6.6FAS set, with all correlation, reduces to

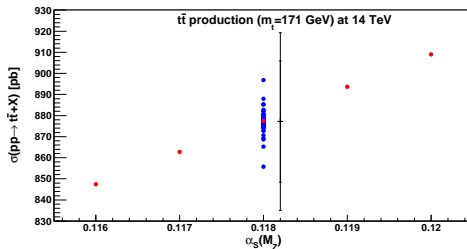
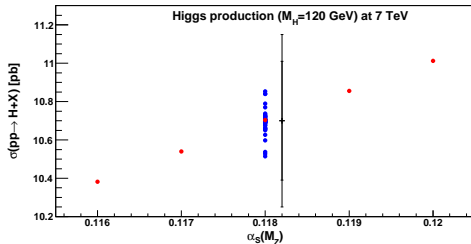
$$\Delta X = \sqrt{\Delta X_{CTEQ6.6}^2 + \Delta X_{\alpha_s}^2},$$

where

- $\Delta X_{CTEQ6.6}$ is the CTEQ6.6 PDF uncertainty from 44 PDFs with the same $\alpha_s(M_Z) = 0.118$
- $\Delta X_{\alpha_s} = (X_{0.120} - X_{0.116})/2$ is the α_s uncertainty computed with two central CTEQ6.6AS PDFs for $\alpha_s(M_Z) = 0.116$ and 0.120

The full proof is given in the paper and backup slides

PDF and α_s uncertainties for $gg \rightarrow H$ and $t\bar{t}$ production



Full and reduced fits with variable α_s : cross sections

Process	CTEQ6.6+CTEQ6.6AS				CTEQ6.6FAS
$t\bar{t}$ (171 GeV)	σ_0	$\Delta\sigma_{PDF}$	$\Delta\sigma_{\alpha_s}$	$\Delta\sigma$	$\sigma_0 \pm \Delta\sigma$
LHC 7 TeV	157.41	10.97	7.54	13.31	160.10 ± 13.93
LHC 10 TeV	396.50	18.75	16.10	24.71	400.48 ± 25.74
LHC 14 TeV	877.19	28.79	30.78	42.15	881.62 ± 44.27
$gg \rightarrow H$ (120 GeV)	σ_0	$\Delta\sigma_{PDF}$	$\Delta\sigma_{\alpha_s}$	$\Delta\sigma$	$\sigma_0 \pm \Delta\sigma$
Tevatron 1.96 TeV	0.63	0.042	0.032	0.053	0.64 ± 0.055
LHC 7 TeV	10.70	0.31	0.32	0.45	10.70 ± 0.48
LHC 10 TeV	20.33	0.66	0.56	0.87	20.28 ± 0.93
LHC 14 TeV	35.75	1.31	0.94	1.61	35.63 ± 1.70
$gg \rightarrow H$ (160 GeV)	σ_0	$\Delta\sigma_{PDF}$	$\Delta\sigma_{\alpha_s}$	$\Delta\sigma$	$\sigma_0 \pm \Delta\sigma$
Tevatron 1.96 TeV	0.26	0.026	0.015	0.030	0.26 ± 0.031
LHC 7 TeV	5.86	0.16	0.18	0.24	5.88 ± 0.26
LHC 10 TeV	11.73	0.33	0.33	0.47	11.72 ± 0.50
LHC 14 TeV	21.48	0.68	0.56	0.88	21.43 ± 0.94
$gg \rightarrow H$ (250 GeV)	σ_0	$\Delta\sigma_{PDF}$	$\Delta\sigma_{\alpha_s}$	$\Delta\sigma$	$\sigma_0 \pm \Delta\sigma$
Tevatron 1.96 TeV	0.055	0.0099	0.0044	0.011	0.058 ± 0.012
LHC 7 TeV	2.30	0.085	0.081	0.12	2.32 ± 0.12
LHC 10 TeV	5.08	0.14	0.15	0.21	5.10 ± 0.22
LHC 14 TeV	10.03	0.26	0.27	0.37	10.04 ± 0.41

The full and reduced methods perfectly agree

Summary

CT09MC PDFs

- a reasonably successful attempt to construct PDFs for LO showering programs
- by their design, the PDFs reproduce key LHC inclusive cross sections at NLO
- compatible with existing tunes of MC programs

Summary II

CTEQ6.6AS PDF sets:

- 4 alternative CTEQ6.6 fits for

$$\alpha_s(M_Z) = 0.116, 0.117, 0.119, 0.120$$

- sufficient to compute uncertainty in $\alpha_s(M_Z)$ at $\approx 68\%$ and 90% C. L., including **the world-average** $\alpha_s(M_Z) = 0.118 \pm 0.002$ as an **input data point**
- **The CTEQ6.6AS** α_s uncertainty should be combined with the CTEQ6.6 PDF uncertainty as

$$\Delta X = \sqrt{\Delta X_{CTEQ6.6}^2 + \Delta X_{CTEQ6.6AS}^2}$$

- The total uncertainty ΔX reproduces the full correlation between $\alpha_s(M_Z)$ and PDFs

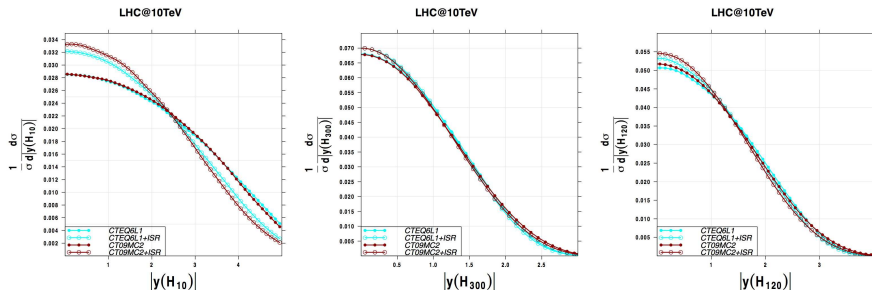
Summary III

- Heavy-quark benchmarking, new statistical methods, and other developments

Backup slides

Effect of parton showering on CT09MC predictions

PYTHIA predictions for $gg \rightarrow H$, with and without contributions from the initial-state radiation

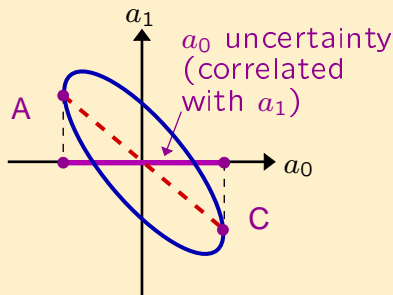
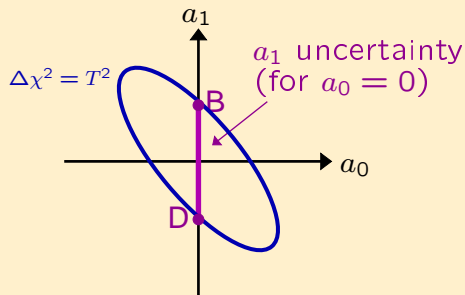


Kinematic showering effects are mild, compared to the PDF dependence, for invariant masses $\gtrsim 100$ GeV

Quadrature theorem for 2 parameters

Physical basis a_i

$$\Delta\chi^2 = \sum_{i,j} H_{i,j} a_i a_j$$

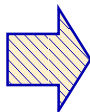
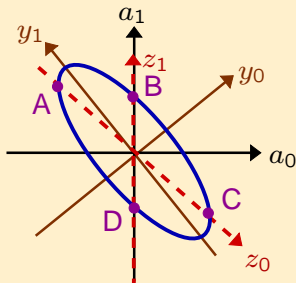


$$\Delta X_1^2 = \frac{1}{4} (X(B) - X(D))^2$$

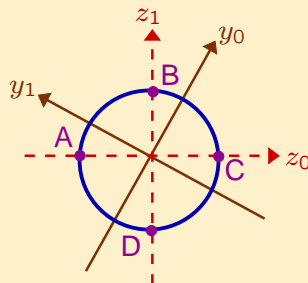
$$\Delta X_0^2 = \frac{1}{4} (X(A) - X(C))^2$$

Quadrature theorem for 2 parameters, cont.

Physical basis a_i
 $\Delta\chi^2 = \sum_{i,j} H_{i,j} a_i a_j$



Eigenvector bases y_i, z_i
 $\Delta\chi^2 = \sum_i y_i^2 = \sum_i z_i^2$

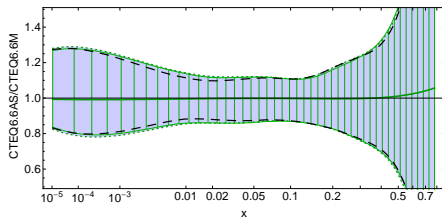


$$\begin{aligned} \Delta X^2 &= \frac{1}{4} \left[(X(A) - X(C))^2 + (X(B) - X(D))^2 \right] \\ &= \Delta X_0^2 + \Delta X_1^2 \end{aligned}$$

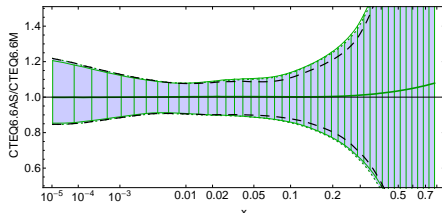
Our findings

Total PDF+ α_s errors ΔX are the **same** when found (a) from a full fit with floating α_s , or (b) by adding ΔX_{PDF} and ΔX_{α_s} in quadrature

g at Q=2 GeV



c at Q=2 GeV



■ black – CTEQ6.6 PDF uncertainty

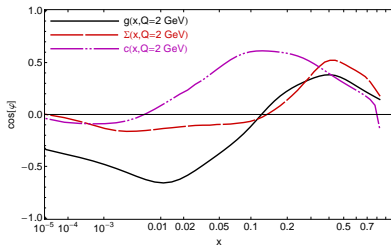
■ Blue filled – PDF+ α_s uncertainty of the fit with floating $\alpha_s(M_Z)$

■ Green hatched – PDF+ α_s uncertainty added in quadrature

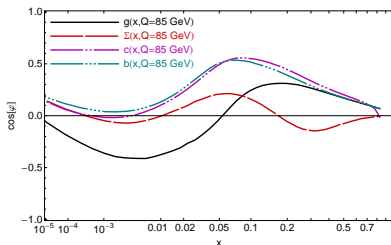
Correlation cosine between CTEQ6.6FAS PDFs and $\alpha_s(M_Z)$

Based on the method in the CTEQ6.6 paper, PRD 78, 013004 (2008)

CTEQ6.6FAS: correlation of $\alpha_s(M_Z)$ with $f_g(x, Q)$



CTEQ6.6FAS: correlation of $\alpha_s(M_Z)$ with $f_g(x, Q)$



$$\cos \varphi = \frac{1}{4\Delta X \Delta Y} \times \sum_{i=1}^{23} \left[\left(X_i^{(+)} - X_i^{(-)} \right) \left(Y_i^{(+)} - Y_i^{(-)} \right) \right]$$

Variations in $\alpha_s(M_Z)$ mostly affect:

- $g(x, Q)$ at $x \approx 0.01$
- $c(x, Q)$ at $x \approx 0.1$
- Singlet quark at $x \approx 0.5$